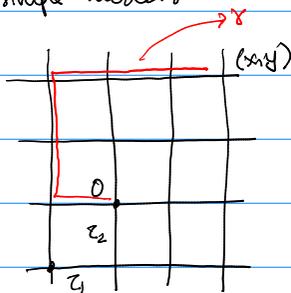


## First Passage Percolation

Simple model.



$$T(x) = \sum_{e \in \gamma} \tau_e$$

$$T(x,y) = \inf_{\gamma: x \rightarrow y} T(\gamma)$$

↓  
first passage time

Extend this to  $\mathbb{R}^d$ .

$$T: \mathbb{R}^d \times \mathbb{R}^d \rightarrow \mathbb{R} \quad \text{in the obvious way.}$$

For any  $x \in \mathbb{R}^d$   $[x]$  is the unique vertex in  $\mathbb{Z}^d$   
 st  $[x] \in [x, x+1)^d$

Limit shape:

$$B(t) = \{y \in \mathbb{R}^d : \tilde{T}(0,y) \leq t\} \quad (\text{random subset of } \mathbb{R}^d)$$



is the passage time smaller than?

It has some sort of growth

Let  $F$  be the cdf of  $\tau_e$

$$\mathbb{P}(\tau_e \leq x) = F(x).$$

$E(\mathbb{Z}^d) = \{ \text{collection of nearest neighbor edges} \}$

$\{\tau_e\}_{e \in E(\mathbb{Z}^d)} := \Omega$  is an iid set of

weights.

OR vertex weights

$$\{\tau_x\}_{x \in \mathbb{Z}^d} := \Omega$$

$$T: \mathbb{Z}^d \times \mathbb{Z}^d \times \Omega \rightarrow \mathbb{R}$$

$$\tilde{T}: \mathbb{R}^d \times \mathbb{R}^d \times \Omega \rightarrow \mathbb{R}$$

$$\tilde{T}(x,y) = T([x], [y])$$

↓  
nearest lattice point to  $x$

Describe the growth of this shape.

Royster  
Durrett

If  $F(0) = 0$  then

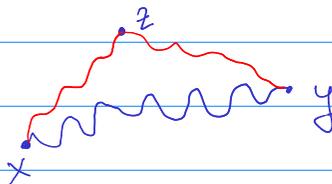
$$P(\tau_e \leq 0) = F(0) \quad (\text{weights are nonnegative})$$

$$\text{and } P(\tau_e = 0) = 0$$

Exer:  $T(x, y)$  is a metric (a.s.)

$$1) T(x, x) = 0$$

$$2) T(x, y) \leq T(x, z) + T(z, y)$$



Typical question:

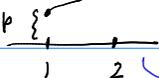
How does  $B(t)$  behave as  $t \rightarrow \infty$ ?

(how does the "metric ball" behave)

$$1) \tau_e \in \text{Uniform}[1, 2]$$

and

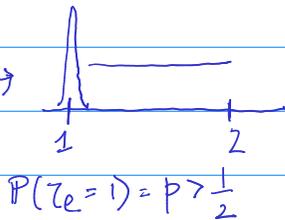
2)  $\tau_e$  has the following cdf



$$P(\tau_e \leq x) = (x-1) \cdot \mathbb{1}_{x \in [1, 2]}$$

see the following videos on the growth of the

ball. (will see (I) first)

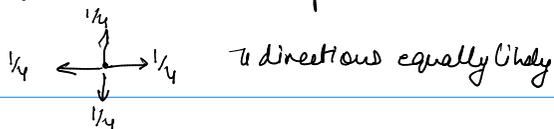


Guess: To first order

$$B(t) = \left\{ (x, y) \in \mathbb{R}^2 : \sqrt{|x|^2 + |y|^2} \leq 1 \right\} t$$

Argument for circular limit shape: take SRW

in 2D,



$X_n$  = displacement of walk.

Pick  $A$ , nice open set.

$$P\left(\frac{X_n}{\sqrt{n}} \in A\right) \xrightarrow{\text{CLT}} \iint_A \frac{e^{-\frac{x^2+y^2}{2}}}{2\pi} dx dy$$

Notice how the limit on the right forgets

the structure of the lattice and acquires

circular symmetry? ( $x^2+y^2$  in the exponent)

(I heard Goussier mention this at some point)

This is REMARKABLE!

So we should expect something similar here.

$$\{\vec{z}_n\}_{n=1}^{\infty} = \text{steps of the walk.}$$

$$\vec{z}_n \in \{\pm e_1, \pm e_2\}$$

$$X_n = \vec{z}_1 + \vec{z}_2 + \dots + \vec{z}_n$$

Normal

density has circular symmetry.

Circular symmetry is "emergent behavior"

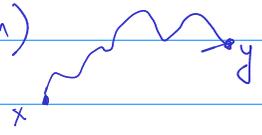
More refined questions

- related → 1) How does  $T(x,y)$  behave as  $|x-y| \rightarrow \infty$
- 2) Does  $B(t)$  have a limit?
- 3) How do geodesics behave?
- 4) How does the cdf of  $\tau_e$  affect all of the above?

5) Why do we care (broader context)  
 - We need basic theorems before giving this  
 and so I'll do this at the end of class.

Question?  $B(t) = t B_0 + o(t)$   
 with high probability.

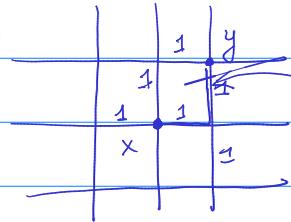
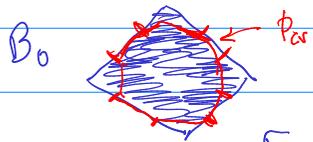
This is a metric. Between far away points (one can show there is an optimal path)



$\tau_e \in \{1, 2\}$  Can you say something about  $B_0$  (assuming it exists)

$$P(\tau_e = 1) = p \in (0, 1).$$

If  $p=1$



'1' path or shortest possible path.

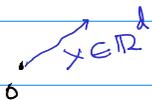
$$T(x,y) = |x-y|,$$

time constant will describe  $\mathbb{B}_0$

Time constant (first order asymptotic)

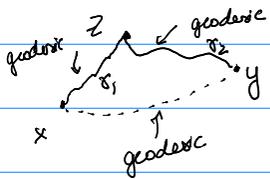
want to understand "average time to go in a particular direction"

$$\dot{N}(t) = Nx$$



$$\frac{T(0, Nx)}{N} = \text{"average time to go in direction x"}$$

Recall  $\Delta$  inequality



$$\begin{aligned} T(x, z) + T(z, y) \\ &= T(\gamma_1) + T(\gamma_2) \\ &= T(\gamma_1 \circ \gamma_2) \\ &\quad \uparrow \\ &\text{path concatenation} \end{aligned}$$

$$\geq T(x, y)$$

$$\text{So } T(0, (N+1)x) \leq T(0, Nx) + T(Nx, (N+1)x)$$

Take expectation

$$E[T(0, (N+1)x)] \leq E[T(0, Nx)] + E[T(Nx, (N+1)x)]$$

$$\mathbb{B}(t) \approx t \mathbb{B}_0$$

But  $\mathbb{B}(t)$  is a subset.

Need some function  $u(x, t)$  st

$\mathbb{B}(t)$  is a subset set

$$\mathbb{B}(t) := \{x \in \mathbb{R}^d : u(x, t) \leq 1\}$$

$$\frac{T(0, Nx)}{N} \xrightarrow{?} \mu_0(x)$$

$$\mathbb{B}_0 = \{x : \mu_0(x) \leq 1\}$$

Subadditivity in  $N_0$

→ triangle inequality

$$= E[T(0, Nx)] + E[T(0, Mx)]$$

important!

$$a_k = E[T(0, kx)]$$

$$a_{n+m} \leq a_n + a_m$$

$\{a_k\}_{k=1}^{\infty}$  is a subadditive sequence

Translation invariance of the measure!

Shift the origin to  $[Nx]$ .

Let us at least state this last result precisely.

$$\Omega = \{T_e\}^{\mathbb{Z}^d}, \quad \mathcal{F} = \text{product } \sigma \text{ algebra}$$

$$\mathbb{P} = \mathbb{P}^{\otimes \mathbb{Z}^d}$$

$\mathbb{P} \leftarrow$  product measure of edge wts.

lemma if  $f: \Omega \rightarrow \mathbb{R}$   $E[|f|] < \infty$

Then if  $f(x+\cdot) := f(\{x+z\}_{z \in \mathbb{Z}^d})$

We must have

$$E[f(x+\cdot)] = E[f(0+\cdot)]$$

Pf: Easy for finitely many variables.

How to do for  $\infty$  variables?

So we have a sequence st

$$a_{n+m} \leq a_n + a_m \quad \forall n, m.$$

Called SUBADDITIVE!

Fekete's Lemma: Let  $\{a_n\}$  be a subadditive sequence.

Then  $\lim_{n \rightarrow \infty} \frac{a_n}{n} = \inf_n \frac{a_n}{n} \in [-\infty, \infty]$

In our case  $0 \leq \frac{a_n}{n} \leq E[\tau_e]$  ← average weight  
EX:  $\tau_e \in \{1, 2\}$   $P(\tau_e = 1) = p$   
 $E[\tau_e] = p + 2(1-p) = 2-p$

pf:  $a_{kn} \leq \underbrace{a_n + a_n + \dots}_{k \text{ times}} = ka_n$  ← induction  
 $\Rightarrow \frac{a_{kn}}{kn} \leq \frac{a_n}{n}$  (Hint at monotonicity)  
 $a_2 = a_{1+1} \leq a_1 + a_1$

So in general, fix  $m$  and look at  $n$  large.

We wish to show that  $\frac{a_m}{m}$  is decreasing

$$n = km + r, \text{ where } r \in \{0, \dots, m-1\}$$

$$a_{km+r} \leq ka_m + a_r \quad (\text{using subadditivity})$$

$$\frac{a_{km+r}}{n} \leq \frac{ka_m + a_r}{n} \leq \frac{a_m}{\frac{m}{k} + \frac{r}{n}} + \frac{a_r}{n}$$

mis fixed  
 $a_r \leq \max\{|a_1|, \dots, |a_{m-1}|\}$

$$\lim_{n \rightarrow \infty} \frac{a_n}{n} \leq \frac{a_m}{m} \quad \lim_{m \rightarrow \infty} \frac{a_m}{m} \Rightarrow \lim_{n \rightarrow \infty} \frac{a_n}{n} \leq \lim_{m \rightarrow \infty} \frac{a_m}{m}$$

Early results.

A generalization of classical percolation theory.

(Hammerley and Weldon 1964): for  $x \in \mathbb{R}^d$

$\frac{T(0, nx)}{n} \rightarrow g(x)$  in probability / measure (Excellent paper to read, used ideas similar to Fekete's lemma)

(Cox & Durrett, Kesten, Richardson)  $\frac{B(t)}{t}$  'converges'

$\frac{E[T(0, nx)]}{n} \rightarrow g(x)$   
average out the randomness.

More refined questions:

1) How to show  $\frac{T(0, nx)}{n} \rightarrow g(x)$  a.s.

(Subadditive ergodic theorem)

$$\lim_{n \rightarrow \infty} \mathbb{P} \left( \left| \frac{T(0, nx)}{n} - g(x) \right| > \epsilon \right) = 0$$

$$\mathbb{P} \left( \lim_{n \rightarrow \infty} \frac{T(0, nx)}{n} = g(x) \right) = 1$$

(We will build some background)

2) In the classical prob, we consider

$$S_n = \sum_{i=1}^n X_i, \quad X_i \text{ iid}$$

$$\frac{S_n}{n} \rightarrow \mu \quad (\text{LLN}) \quad \text{a.s.}$$

Weak law  $\frac{S_n}{n} \rightarrow \mu$  in prob.  $E[X_1]$

Strong law  $\frac{S_n}{n} \rightarrow \mu$  a.s.

CLT  $\uparrow$

THEN we show

$$\frac{S_n - n\mu}{\sqrt{n}} \Rightarrow N(0, 1)$$

We will not be able to prove.

for  $\frac{T(0, nx)}{n}$   
 $\downarrow$   
 $g(x)$  will be unknown

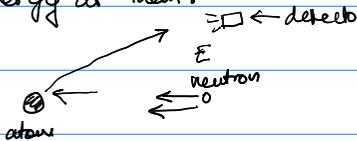
So can we show

$$\frac{T(0, mx) - ng(x)}{\sqrt{n}} \xrightarrow{d} N(0, 1)$$

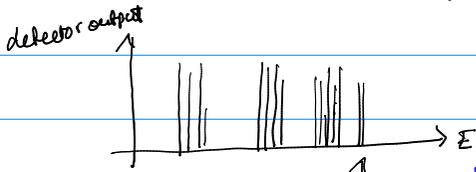
← does not converge to a Normal distribution.

THIS TURNS OUT TO BE FALSE (simuldrn)  
and has an interesting history.

Problem 1: Take an example from Physics. A neutron scattering experiment takes one atom and fires a bunch of neutrons at a given energy at them.



Neutrons have energy  $E$ . At most energies, neutrons pass right through, but at some  $E$ , the neutrons are SCATTERED by the atom.



↑ cannot say where exactly the resonance happens

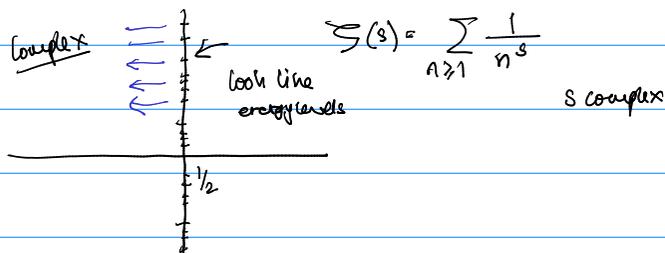
This discreteness is due to QM.

For simple atoms like hydrogen, the picture is simple, and we can predict what these energies look like. (Balmer Lines?)

But for heavy atoms (lead, Thorium, Lanthanide series) it's impossible to solve QM equations

- Many, many dense but discrete resonances
- Best description is as a point process.
- Statistics (gaps, k-point correlation) are UNIVERSAL (for heavy enough atoms)

## Problem 2: Zeta function



What's the distance between zeros of the zeta function?

Same statistics.



Montgomery.

(Zeta  $\zeta$  results under RH)

- RM results universality shown Tao-Vu,

Erdos Schlicien Yau et al (2010-2015)

Original Gaussian results (1950s-1980s)

Wigner, Dyson, Mehta, Tracy, Widom, etc.

Add